# Teleparallel Newton-Cartan gravity 

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| 1 | 1 | Leibniz |
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| 10 | 2 | Universität |
| 100 | 4 | Hannover |

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## Motivation

- 'Non-relativistic' limit: GR $\rightarrow$ Newtonian gravity
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- Galilei-relativistic spacetime geometry
- Newtonian gravity described by curved connection


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- Textbook formulation: linearised gravity
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- Geometric formulation: GR $\xrightarrow{c \rightarrow \infty}$ Newton-Cartan gravity
- Galilei-relativistic spacetime geometry
- Newtonian gravity described by curved connection
- Geometric description of the $c \rightarrow \infty$ limit of teleparallel gravity?
- Here: consider just TEGR
- Read, Teh (2018): special case of our formalism; null-reduction instead of $c \rightarrow \infty$

J Read, NJ Teh: The teleparallel equivalent of Newton-Cartan gravity, arXiv:1807.11779, Class. Quantum Gravity 35, 18LT01 (2018)

## Outline

(1) Recap of usual Newton-Cartan gravity

- Galilei manifolds
- Newton-Cartan gravity
(2) Teleparallel Galilei connections
- Bargmann structures
- Teleparallel Newton-Cartan gravity

3 Relation to other theories

- Teleparallel NC from TEGR
- Recovering Newtonian gravity


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PKS: Teleparallel Newton-Cartan gravity, in preparation

## Galilei manifolds

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- $\int_{\gamma} \tau=$ elapsed time along $\gamma, h$ defines metric on $\operatorname{ker} \tau$


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- $\int_{\gamma} \tau=$ elapsed time along $\gamma, h$ defines metric on $\operatorname{ker} \tau$
- This talk: mostly assume $\mathrm{d} \tau=0$ (absolute time)


## Galilei frames

- Homogeneous Galilei group:

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\begin{equation*}
\mathrm{Gal}=\mathrm{O}(3) \ltimes \mathbb{R}^{3} \tag{1}
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- Galilei frame for $(M, \tau, h)$ : local frame $\left(\mathrm{e}_{A}\right)=\left(\mathrm{e}_{t}=v, \mathrm{e}_{a}\right)$ s.t. $\tau(v)=1$, $h^{\mu v}=\delta^{a b} \mathrm{e}_{a}^{\mu} \mathrm{e}_{b}^{v}$


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- $\rightsquigarrow$ Galilei frame bundle $G(M) \xrightarrow{\pi} M$, principal Gal-bundle


## Galilei connections

- Galilei connection: connection $\omega$ on $G(M)$
- $\nabla \tau=0=\nabla h$, or $\left(\omega^{a}{ }_{b}, \omega^{a}\right)$ valued in $\mathfrak{g a l}=\mathfrak{s o}(3) \oplus \mathbb{R}^{3}$


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- $\Longrightarrow$ Determined by torsion $\left(T^{t}, T^{a}\right)=\left(\mathrm{d} \tau, T^{a}\right)$ and Newton-Coriolis form

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- Newtonian connection: $T=0$ and $\mathrm{d} \Omega=0$


## Newton-Cartan gravity

## Axioms for Newton-Cartan gravity

- Spacetime is a Galilei manifold $(M, \tau, h)$ with absolute time, endowed with a Newtonian connection $\widetilde{\omega}$,
(2) ideal clocks measure time as defined by $\tau$,
(3) free test particles move on timelike geodesics of $\widetilde{\omega}$, i.e. timelike curves $\gamma$ solving

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holds, where $\rho$ is the mass density.

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- Arises as formal $c \rightarrow \infty$ limit of GR
- Coordinate formulation of Newtonian gravity can be recovered (up to possibility of non-absolute rotation)


## Teleparallel Galilei connections

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## 'Teleparallelisation' of NC gravity?

- 'Teleparallelisation' GR $\rightarrow$ TEGR:
- Lorentzian manifold $(M, g)$ : unique torsion-free metric connection $\stackrel{\llcorner }{\boldsymbol{\omega}}$
- $\Longrightarrow$ General metric connection $\stackrel{L}{\omega}$ determined by its torsion $\stackrel{\llcorner }{T}$
- $\rightsquigarrow$ Riemannian curvature $\stackrel{\llcorner }{\tilde{R}}$ expressible purely in terms of $\stackrel{\llcorner }{\omega}, \stackrel{\llcorner }{T}, \stackrel{\llcorner }{R}=0$
- $\rightsquigarrow$ Reformulation of GR in terms of $\stackrel{L}{\omega}, \stackrel{L}{T}$


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- $\rightsquigarrow$ Reformulation of GR in terms of $\stackrel{\llcorner }{\omega}, \stackrel{\llcorner }{T}$
- Analogue for NC gravity?
- Problem: Galilei connection $\omega$ not uniquely determined by $T$ !


## Extending Galilei to Bargmann structures

- Bargmann group Barg $=\mathrm{Gal} \ltimes\left(\mathbb{R}^{4} \times \mathrm{U}(1)\right)$
- Extend Galilei frame bundle to Barg-bundle $B(M)=G(M) \times{ }_{\text {Gal }}$ Barg


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- Connection on $B(M)=$ Galilei connection $\boldsymbol{\omega}+$ tensorial form $(\boldsymbol{\theta}, \mathrm{i} \boldsymbol{a}) \in \Omega_{\mathrm{Gal}}^{1}\left(G(M), \mathbb{R}^{4} \oplus \mathfrak{u}(1)\right)$
'Globalised' construction from:
M Geracie, K Prabhu, MM Roberts: Curved non-relativistic spacetimes, Newtonian gravitation and massive matter, arXiv:1503.02682, J. Math. Phys. 56, 103505 (2015)


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- Bargmann structure: Choice of $\boldsymbol{a}$ with $\boldsymbol{\theta}$ corresponding to canonical solder form
- Pulled back to $M$ : extended coframe $\left(\mathrm{e}^{t}=\tau, \mathrm{e}^{a}, \mathrm{i} a\right)$
- Transformation under local boost of frame $v \rightarrow v-k^{a} \mathrm{e}_{a}$ :

$$
\begin{equation*}
a \rightarrow a+k_{a} \mathrm{e}^{a}+\frac{1}{2}|k|^{2} \tau \tag{6}
\end{equation*}
$$

## Extended torsion

- Exterior covariant derivative of $\left(\tau, \mathrm{e}^{a}, \mathrm{i} a\right)$ : extended torsion

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\begin{equation*}
\mathrm{d}^{\omega}\left(\tau, \mathrm{e}^{a}, \mathrm{i} a\right)=\left(T^{t}, T^{a}, \mathrm{i} f\right) \tag{7}
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- Mass torsion $f=\mathrm{d} a+\omega_{a} \wedge \mathrm{e}^{a}=\mathrm{d} a+\Omega$


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- For general $\boldsymbol{\omega}$ : Newton-Cartan contortion

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\begin{equation*}
\Gamma_{\mu \nu}^{\rho}-\widetilde{\Gamma}_{\mu \nu}^{\rho}=\frac{1}{2} T_{\mu \nu}^{\rho}-T_{(\mu v)}^{\rho}+\tau_{(\mu} f_{v)}^{\rho}=: K_{\mu v}^{\rho} \tag{8}
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## Teleparallel Newton-Cartan gravity

## Axioms for teleparallel Newton-Cartan gravity

(1) Spacetime is a Galilei manifold $(M, \tau, h)$ with absolute time, endowed with a Bargmann structure and a flat Galilei connection $\omega$,
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- the field equation

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\begin{equation*}
-D_{\sigma} K^{\sigma}{ }_{A B}+D_{A} K_{\mu B}^{\mu}-K_{\sigma B}^{\mu} T^{\sigma}{ }_{\mu A}+K_{\mu \sigma}^{\mu} K_{A B}^{\sigma}-K_{A \sigma}^{\mu} K^{\sigma}{ }_{\mu B}=4 \pi G \rho \tau_{A} \tau_{B} \tag{10}
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- LHS of (10) is $\widetilde{R}_{A B} \rightsquigarrow$ equivalent to usual NC gravity


## Relation to other theories

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- Lorentzian tetrad (assume $\mathrm{d} \tau=0$ ):

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- Lorentzian local connection form:

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- Field equation of TEGR $\rightsquigarrow$ trace-reverse $\xrightarrow{c \rightarrow \infty}$ teleparallel NC field eq.!


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- For $\left[\mathrm{e}_{a}, v\right]=0$ : equation of motion becomes

$$
\begin{equation*}
\ddot{\gamma}^{a}+\omega_{c}^{a}{ }_{b} \dot{\gamma}^{c} \dot{\gamma}^{b}=-\partial^{a} \phi \tag{14}
\end{equation*}
$$

## Conclusion

## Summary

- Bargmann structure formalism for NC geometry $\rightsquigarrow$ natural notion of teleparallel Galilei connections
- Teleparallel formulation of NC gravity
- Arises from TEGR for $c \rightarrow \infty$, reproduces Newtonian gravity


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## Outlook

- TNC generalisation $(\mathrm{d} \tau \neq 0)$ ? $\rightsquigarrow$ probably needs different symmetry algebra ('TNC type I')
- Action formulation?
- Modified teleparallel gravity?
- Theories with more general geometry?
- 'Covariant' post-Newtonian limit of (modified) teleparallel gravity?


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## Many thanks for your attention!

## Appendix: details

4 Details on recovery of Newtonian gravity

## Gauge fixing the purely spatial torsion

- Purely spatial part of field equation: $\widetilde{R}_{a b}=0 \stackrel{3 D}{ }$ spatial metric flat
- $\Longrightarrow$ We may assume

$$
\begin{equation*}
T_{b c}^{a}=0 \tag{15}
\end{equation*}
$$

consistently with flatness

## Trautman's 'absolute rotation' condition

- Usual NC: To recover Newtonian gravity, assume $\widetilde{R}^{a b}{ }_{\mu \nu}=0$
$\bullet \Longleftrightarrow \exists$ rigid, non-rotating frames
- Such frames in teleparallel NC:

$$
\begin{align*}
\omega_{(a b)} & =T_{(a b) t}  \tag{16a}\\
\omega_{[a b]} & =\frac{1}{2} f_{a b} \tag{16b}
\end{align*}
$$

